

GaAs - Gallium Arsenide

Electrical properties

[Basic Parameters](#)

[Mobility and Hall Effect](#)

[Transport Properties in High Electric Fields](#)

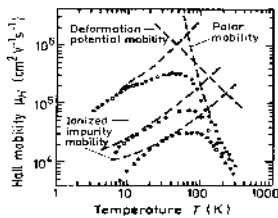
[Impact Ionization](#)

[Recombination Parameters](#)

Basic Parameters

Breakdown field	$\approx 4 \cdot 10^5$ V/cm
Mobility electrons	≤ 8500 cm ² V ⁻¹ s ⁻¹
Mobility holes	≤ 400 cm ² V ⁻¹ s ⁻¹
Diffusion coefficient electrons	≤ 200 cm ² /s
Diffusion coefficient holes	≤ 10 cm ² /s
Electron thermal velocity	$4.4 \cdot 10^5$ m/s
Hole thermal velocity	$1.8 \cdot 10^5$ m/s

Mobility and Hall Effect



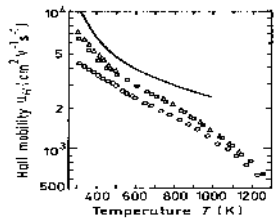
Electron Hall mobility versus temperature for different doping levels.

([Stillman et al. \[1970\]](#))

1. Bottom curve: $N_d = 5 \cdot 10^{15}$ cm⁻³;
2. Middle curve: $N_d = 10^{15}$ cm⁻³;
3. Top curve: $N_d = 5 \cdot 10^{15}$ cm⁻³

For weakly doped GaAs at temperature close to 300 K, electron Hall mobility

$$\mu_H = 9400(300/T) \text{ cm}^2 \text{ V}^{-1} \text{ s}^{-1}$$



Electron Hall mobility versus temperature for different doping levels and degrees of compensation (high temperatures):

Open circles: $N_d = 4N_a = 1.2 \cdot 10^{17}$ cm⁻³;

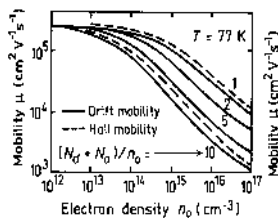
Open squares: $N_d = 4N_a = 10^{16}$ cm⁻³;

Open triangles: $N_d = 3N_a = 2 \cdot 10^{15}$ cm⁻³;

Solid curve represents the calculation for pure GaAs ([Blakemore \[1982\]](#)).

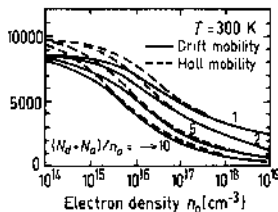
For weakly doped GaAs at temperature close to 300 K, electron drift mobility

$$\mu_n = 8000(300/T)^{2/3} \text{ cm}^2 \text{ V}^{-1} \text{ s}^{-1}$$



Drift and Hall mobility versus electron concentration for different degrees of compensation $T = 77$ K

([Rode \[1975\]](#)).



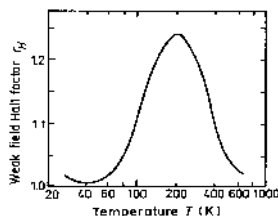
Drift and Hall mobility versus electron concentration for different degrees of compensation $T = 300$ K

([Rode \[1975\]](#)).

Approximate formula for the Hall mobility

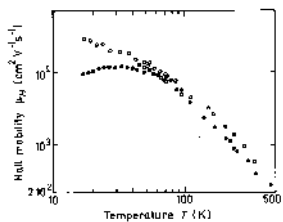
$$\mu_n = \mu_{OH} / (1 + N_d 10^{-17})^{1/2}, \text{ where } \mu_{OH} \approx 9400 \text{ (cm}^2 \text{ V}^{-1} \text{ s}^{-1}), N_d \text{ in cm}^{-3}$$

([Hilsum \[1974\]](#)).



Temperature dependence of the Hall factor for pure n -type GaAs in a weak magnetic field

([Rode \[1975\]](#)).



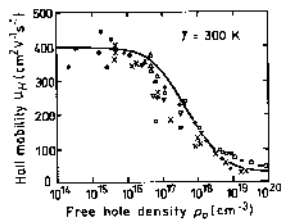
Temperature dependence of the Hall mobility for three high-purity samples
(Wiley [1975]).

For GaAs at temperatures close to 300 K, hole Hall mobility

$$\mu_{pH} = \left[0.0025 \left(\frac{T}{300} \right)^{2.3} + 4 \times 10^{21} p \left(\frac{T}{300} \right)^{1.5} \right]^{-1} \text{ (cm}^2 \text{V}^{-1} \text{s}^{-1}), \text{ (p - in cm}^{-3}\text{)}$$

For weakly doped GaAs at temperature close to 300 K, Hall mobility

$$\mu_{pH} = 400(300/T)^{2.3} \text{ (cm}^2 \text{V}^{-1} \text{s}^{-1}\text{)}.$$

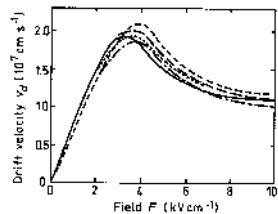


The hole Hall mobility versus hole density.
(Wiley [1975]).

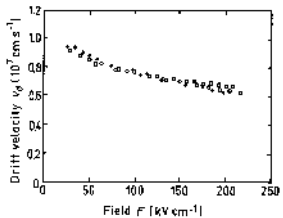
At $T = 300$ K, the Hall factor in pure GaAs

$$r_H = 1.25.$$

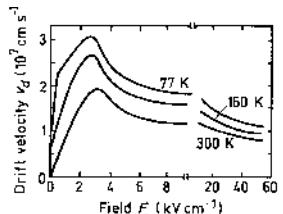
Transport Properties in High Electric Fields



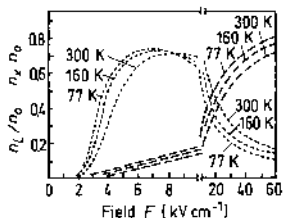
Field dependences of the electron drift velocity.
(Blakemore [1982]).
Solid curve was calculated by (Pozhela and Reklaitis [1980]).
Dashed and dotted curves are measured data, 300 K



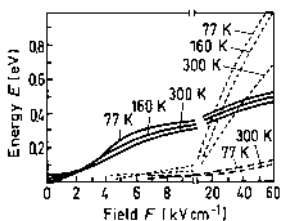
Field dependences of the electron drift velocity for high electric fields, 300 K.
(Blakemore [1982]).



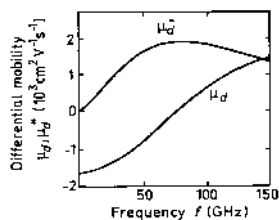
Field dependences of the electron drift velocity at different temperatures.
(Pozhela and Reklaitis [1980]).



Fraction of electrons in L and X valleys. n_L and n_X as a function of electric field F at 77, 160, and 300 K, $N_d = 0$
(Pozhela and Reklaitis [1980]).
Dotted curve - L valleys, dashed curve - X valleys.



Mean energy E in Γ , L , and X valleys as a function of electric field F at 77, 160, and 300 K, $N_d = 0$
(Pozhela and Reklaitis [1980]).
Solid curve - Γ valleys, dotted curve - L valleys, dashed curve - X valleys.

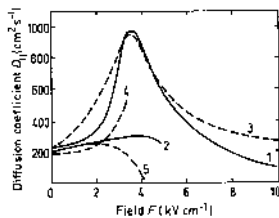


Frequency dependences of electron differential mobility.

μ_d is real part of the differential mobility; μ_d^* is imaginary part of differential mobility.

$F = 5.5 \text{ kV cm}^{-1}$

(Rees[1969]).



The field dependence of longitudinal electron diffusion coefficient $D//F$.

Solid curves 1 and 2 are theoretical calculations. Dashed curves 3, 4, and 5 are experimental data.

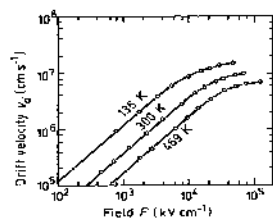
Curve 1 - from (Pozhela and Reklaitis[1980]).

Curve 2 - from (Fauquembergue et al.[1980]).

Curve 3 - from (Ruch and Kinof[1968]).

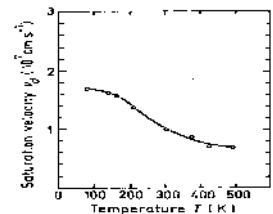
Curve 4 - from (Bareikis et al.[1978]).

Curve 5 - (from de Murcia[1991]).



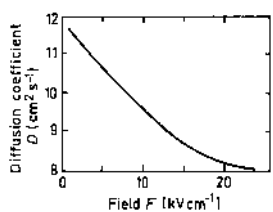
Field dependences of the hole drift velocity at different temperatures.

(Datal et al. [1971]).



Temperature dependence of the saturation hole velocity in high electric fields

(Datal et al. [1971]).



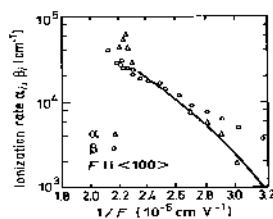
The field dependence of the hole diffusion coefficient.

(Joshi and Crendin [1989]).

Impact Ionization

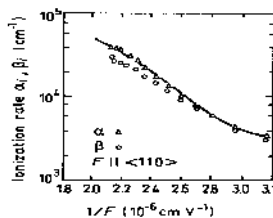
There are two schools of thought regarding the impact ionization in GaAs.

The first one states that impact ionization rates α_i and β_i for electrons and holes in GaAs are known accurately enough to distinguish such subtle details such as the anisotropy of α_i and β_i for different crystallographic directions. This approach is described in detail in the work by Dmitriev et al.[1987].



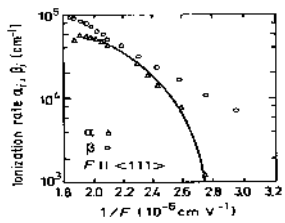
Experimental curves α_i and β_i versus $1/F$ for GaAs.

(Pearsall et al. [1978]).



Experimental curves α_i and β_i versus $1/F$ for GaAs.

(Pearsall et al. [1978]).



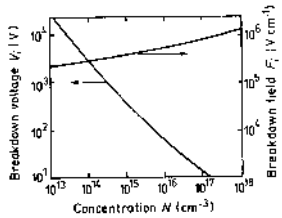
Experimental curves α_i and β_i versus $1/F$ for GaAs.

([Pearsall et al. \[1978\]](#)).

The second school focuses on the values of α_i and β_i for the same electric field reported by different researches differ by an order of magnitude or more. This point of view is explained by Kyuregyan and Yurkov [1989]. According to this approach we can assume that $\alpha_i = \beta_i$. Approximate formula for the field dependence of ionization rates:

$$\alpha_i = \beta_i = \alpha_0 \exp[\delta - (\delta^2 + (F_0/F)^2)^{1/2}]$$

where $\alpha_0 = 0.245 \cdot 10^6 \text{ cm}^{-1}$; $\beta = 57.6 F_0 = 6.65 \cdot 10^6 \text{ V cm}^{-1}$ (Kyuregyan and Yurkov [1989]).



Breakdown voltage and breakdown field versus doping density for an abrupt p - n junction.

([Kyuregyan and Yurkov \[1989\]](#)).

Recombination Parameter

Pure n-type material ($n_0 \sim 10^{14} \text{ cm}^{-3}$)

The longest lifetime of holes $\tau_p \sim 3 \cdot 10^{-6} \text{ s}$

Diffusion length $L_p = (D_p \tau_p)^{1/2}$ $L_p \sim 30\text{-}50 \text{ }\mu\text{m}$.

Pure p-type material

(a) Low injection level

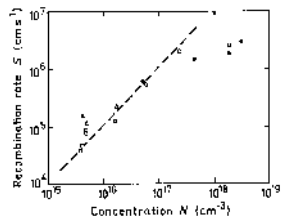
The longest lifetime of electrons $\tau_n \sim 5 \cdot 10^{-9} \text{ s}$

Diffusion length $L_n = (D_n \tau_n)^{1/2}$ $L_n \sim 10 \text{ }\mu\text{m}$

(b) High injection level (filled traps)

The longest lifetime of electrons $\tau \sim 2.5 \cdot 10^{-7} \text{ s}$

Diffusion length L_n $L_n \sim 70 \text{ }\mu\text{m}$



Surface recombination velocity versus doping density

([Aspnes \[1983\]](#)).

Different experimental points correspond to different surface treatment methods.

Radiative recombination coefficient ([Varshni \[1967\]](#))

90 K $1.8 \cdot 10^{-8} \text{ cm}^3/\text{s}$

185 K $1.9 \cdot 10^{-9} \text{ cm}^3/\text{s}$

300 K $7.2 \cdot 10^{-10} \text{ cm}^3/\text{s}$

Auger coefficient

300 K $\sim 10^{-30} \text{ cm}^6/\text{s}$

500 K $\sim 10^{-29} \text{ cm}^6/\text{s}$

